

Robot Dynamics

Dynamics

• For Manipulators, the general form is

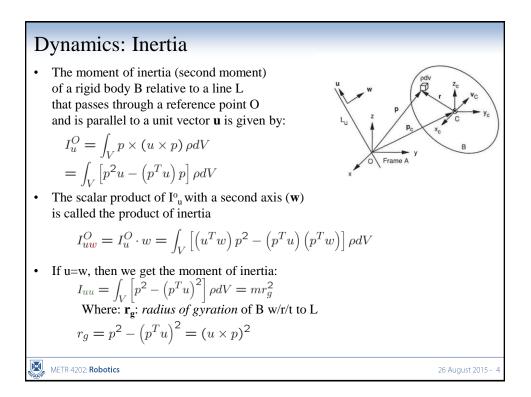
$$\tau = M(\Theta)\ddot{\Theta} + V(\Theta, \dot{\Theta}) + G(\Theta)$$

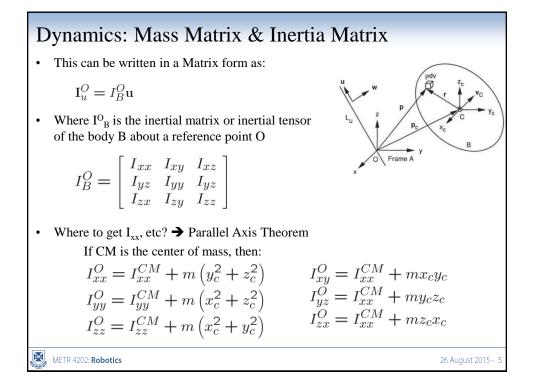
where

- τ is a vector of joint torques
- Θ is the nx1 vector of joint angles
- $M(\Theta)$ is the nxn mass matrix
- $V(\Theta, \Theta)$ is the nx1 vector of centrifugal and Coriolis terms
- $G(\Theta)$ is an nx1 vector of gravity terms
- Notice that all of these terms depend on Θ so the dynamics varies as the manipulator move

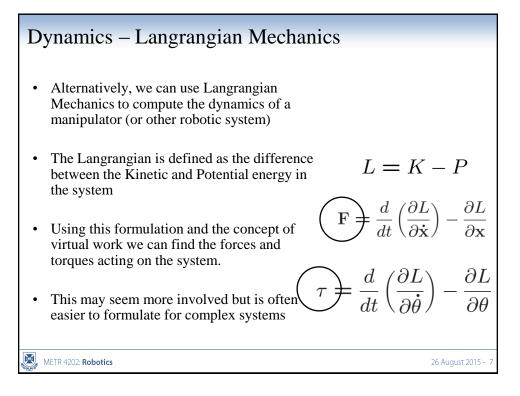
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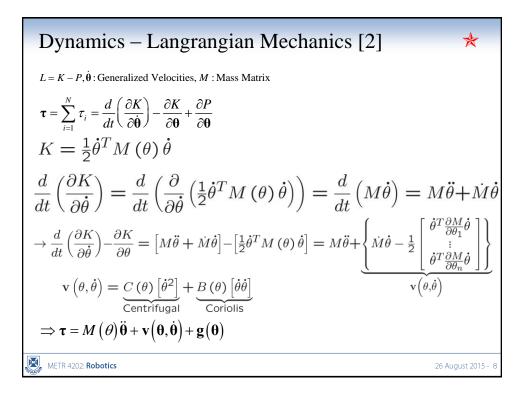
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Dynamics: Mass Matrix • The Mass Matrix: Determining via the Jacobian! $\kappa = \sum_{i=1}^{N} \kappa_{i}$ $\kappa_{i} = \frac{1}{2} \left(m_{i} v_{C_{i}}^{T} v_{C_{i}} + \omega_{i}^{T} I_{C_{i}} \omega_{i} \right)$ $v_{C_{i}} = J_{v_{i}} \dot{\theta} \quad J_{v_{i}} = \begin{bmatrix} \frac{\partial p_{C_{1}}}{\partial \theta_{1}} & \cdots & \frac{\partial p_{C_{i}}}{\partial \theta_{i}} & \underbrace{0}_{i+1} & \cdots & \underbrace{0}_{n} \end{bmatrix}$ $\omega_{i} = J_{\omega_{i}} \dot{\theta} \quad J_{\omega_{i}} = \begin{bmatrix} \overline{\varepsilon}_{1} Z_{1} & \cdots & \overline{\varepsilon}_{i} Z_{i} & \underbrace{0}_{i+1} & \cdots & \underbrace{0}_{n} \end{bmatrix}$ $\therefore M = \sum_{i=1}^{N} \left(m_{i} J_{v_{i}}^{T} J_{v_{i}} + J_{\omega_{i}}^{T} I_{C_{i}} J_{\omega_{i}} \right)$! M is symmetric, positive definite $\therefore m_{ij} = m_{ji}, \dot{\theta}^{T} M \dot{\theta} > 0$

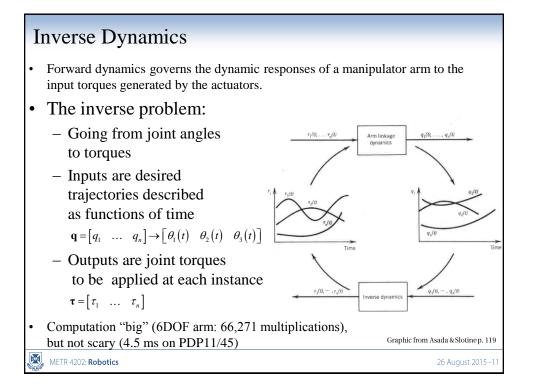




Dynamics – Langrangian Mechanics [3]
• The Mass Matrix: Determining via the Jacobian!

$$\begin{split} & \kappa = \sum_{i=1}^{N} \kappa_i \\ & K_i = \frac{1}{2} \left(m_i v_{C_i}^T v_{C_i} + \omega_i^T I_{C_i} \omega_i \right) \\ & v_{C_i} = J_{v_i} \dot{\theta} \quad J_{v_i} = \begin{bmatrix} \frac{\partial \mathbf{p}_{C_1}}{\partial \theta_1} & \cdots & \frac{\partial \mathbf{p}_{C_i}}{\partial \theta_i} & \underbrace{\mathbf{0}}_{i+1} & \cdots & \underbrace{\mathbf{0}}_{n} \end{bmatrix} \\ & \omega_i = J_{\omega_i} \dot{\theta} \quad J_{\omega_i} = \begin{bmatrix} \overline{\varepsilon}_1 Z_1 & \cdots & \overline{\varepsilon}_i Z_i & \underbrace{\mathbf{0}}_{i+1} & \cdots & \underbrace{\mathbf{0}}_{n} \end{bmatrix} \\ & \therefore M = \sum_{i=1}^{N} \left(m_i J_{v_i}^T J_{v_i} + J_{\omega_i}^T I_{C_i} J_{\omega_i} \right) \\ & ! \text{ M is symmetric, positive definite } \therefore m_{ij} = m_{ji}, \dot{\mathbf{\theta}}^T M \dot{\mathbf{\theta}} > 0 \end{split}$$

Generalized Coordinates A significant feature of the Lagrangian Formulation is that any convenient coordinates can be used to derive the system. Go from Joint → Generalized Define p: dp = Jdq q = [q₁ ... q_n] → p = [p₁ ... p_n] Thus: the kinetic energy and gravity terms become KE = ½ ṗ^TH*ṗ G* = (J⁻¹)^TG where: H* = (J⁻¹)^T HJ⁻¹



Also: Inverse Jacobian

• In many instances, we are also interested in computing the set of joint velocities that will yield a particular velocity at the end effector

$$\dot{\theta} = \mathbf{J}(\theta)^{-1} \dot{\mathbf{X}}$$

- We must be aware, however, that the inverse of the Jacobian may be undefined or singular. The points in the workspace at which the Jacobian is undefined are the *singularities* of the mechanism.
- Singularities typically occur at the workspace boundaries or at interior points where degrees of freedom are lost

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